

THE FLOW OF A VISCOUS, ELECTRICALLY CONDUCTING GAS IN A TRANSVERSE MAGNETIC FIELD IN THE PRESENCE OF HEAT TRANSFER

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This paper contains a generalization of the solution for plane flow presented in [1], in that it includes the effect of viscous dissipation.

The problem of the two-dimensional flow of an incompressible, viscous, electrically conducting gas through a channel formed by two parallel planes subjected to the action of a uniform magnetic field of intensity H_0 , and in the presence of heat transfer from the walls, was considered in [1,2]; the first reference contained a solution for a uniform heat flux, but excluded viscous dissipation; the second reference contained a solution which included the effect of viscous dissipation, but was restricted to the case of a wall of constant temperature.

We shall consider the flow of a fluid through a channel formed by two infinite, electrically non-conducting planes at $z = \pm b$, exposed to a normal uniform magnetic field H_0 . The properties of the liquid are described by its electric conductivity σ , its density ρ , specific heat c_p , viscosity μ and thermal conductivity λ . The effect of temperature on these properties will be ignored.

If the walls are impermeable to the fluid, it is possible to consider that the streamlines are confined in planes $x - y$. We choose the axis y to make the velocity component $W_y = 0$. Then the solution can be written

$$\begin{aligned} W_x = W(z), & \quad W_y = 0, & \quad W_z = 0, \\ H_x = H_x(z), & \quad H_y = 0, & \quad H_z = H_0, & \quad p = p(x, z), & \quad T = T(x, z) \end{aligned} \quad (1)$$

the remaining parameters being constant.

The system of equations satisfied by solution (1) has the form

$$\begin{aligned} \frac{\partial}{\partial x} \left(p + \mu_e \frac{H^2}{2} \right) &= \mu \frac{d^2 W}{dz^2} + \mu_e H_0 \frac{dH_x}{dz} \quad \left(\alpha = \frac{\lambda}{\rho c_p} \right) \\ \frac{\partial}{\partial z} \left(p + \mu_e \frac{H^2}{2} \right) &= 0, \quad H_0 \frac{dW}{dz} + \frac{1}{\mu_e \sigma} \frac{d^2 H_x}{dz^2} = 0 \\ W \frac{\partial T}{\partial x} &= \alpha \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial z^2} \right) + \frac{1}{\rho c_p \sigma} \left(\frac{dH_x}{dz} \right)^2 + \frac{\mu}{\rho c_p} \left(\frac{dW}{dz} \right)^2 \end{aligned} \tag{2}$$

Here α denotes the coefficient of heat transfer. The first three equations yield the following solution for W and H_x [3]:

$$\begin{aligned} W &= \frac{PN}{\mu_e^2 \sigma H_0^2} \frac{\cosh N - \cosh(Nz/b)}{\sinh N}, \quad H_x = \frac{Pb}{\mu_e H_0} \left[\frac{\sinh(Nz/b)}{\sinh N} - \frac{z}{b} \right] \\ P &= -\partial p / \partial x = \text{const}, \quad N = \mu_e H_0 b \sqrt{\sigma / \mu} \end{aligned}$$

Here N denotes the Hartman number and b is the channel half-width.

In the present problem the distribution of heat sources is independent of x . It follows that the temperature profiles are similar in all cross-sections of the channel, and the temperature $T = \tau x + \theta(z)$. Thus, the fourth equation (2) can be reduced to the form

$$\begin{aligned} &\frac{PN\tau}{\mu_e^2 \sigma H_0^2} \left[\frac{\cosh N - \cosh(Nz/b)}{\sinh N} \right] = \\ &= \alpha \frac{d^2 \theta}{dz^2} + \frac{P^2}{\mu_e^2 \sigma H_0^2 \rho c_p} \left[\frac{N \cosh(Nz/b)}{\sinh N} - 1 \right]^2 + \frac{\mu}{\rho c_p} \left[\frac{PN^2}{\mu_e^2 \sigma H_0^2 b} \frac{\sinh(Nz/b)}{\sinh N} \right]^2 \end{aligned} \tag{3}$$

The value of τ can be found from the heat balance

$$\rho W^0 c_p b \tau = q + \int_0^b \left[\frac{1}{\sigma} \left(\frac{dH_x}{dz} \right)^2 + \mu \left(\frac{dW}{dz} \right)^2 \right] dz \quad \left(W^0 = \frac{P}{\mu_e^2 \sigma H_0^2} (N \coth N - 1) \right) \tag{4}$$

where W^0 denotes the average velocity.

Solving (3) and (4) subject to $d\theta/dz = 0$ for $z = 0$ and $\theta = 0$ for $z = \pm b$, we obtain

$$\begin{aligned} T &= \tau x + \tau \frac{Pb^4}{2\alpha\mu N \sinh N} \left[\left(\frac{z}{b} \right)^2 \cosh N - \frac{2}{N^2} \cosh \left(N \frac{z}{b} \right) - \frac{N^2 - 2}{N^2} \cosh N \right] - \frac{P^2 b^4}{\alpha \rho c_p \mu N^2} \\ &\quad \left\{ \frac{1}{4 \sinh^2 N} \left[\frac{1}{2} \cosh \left(2N \frac{z}{b} \right) - \frac{1}{2} \cosh 2N + N^2 \left(\frac{z^2}{b^2} - 1 \right) \right] - \frac{2 \cosh(Nz/b)}{N \sinh N} \right\} + \\ &+ \frac{2 \cosh N}{N \sinh N} + \frac{1}{2} \left(\frac{z^2}{b^2} - 1 \right) \left\{ - \frac{P^2 b^4 N^2}{4\alpha \rho c_p \mu \sinh^2 N} \left[\frac{1}{2} \cosh \left(2N \frac{z}{b} \right) - \frac{1}{2} \cosh 2N + \left(\frac{z^2}{b^2} - 1 \right) (N \cosh N - \sinh N) \right] \right\} \\ \theta &= \frac{1}{\rho c_p b W^0} \left[q + \frac{P^2 b^3}{\mu} \left(\frac{2N}{2N \sinh^2 N} - \frac{1}{N^2} \right) \right] \end{aligned}$$

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